

# SU(N) gauge theories with symmetric-rep fermions

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with

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Part I.  $SU(N)$ ,  $N = 2, 3, 4$ , with  $N_f = 2$  two-index sym rep Dirac fermions:  
“continuum” issues; physics results

Part II: Fat (nHYP) links and what they do for us  
⇒ It's a good idea to try DWF with fat links

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## Walking [extended] Technicolor

- Confinement & Chiral symmetry breaking are lost for  $N_f > N_{\text{critical}}$
- “Walking” theories are found just below the conformal window
- Condensate enhancement needed  $\Leftrightarrow$  mass anomalous dimension  $\gamma \approx 1$
- Natural candidates:  $N_f = 2$  higher-irrep fermions [Sannino *et al.*]

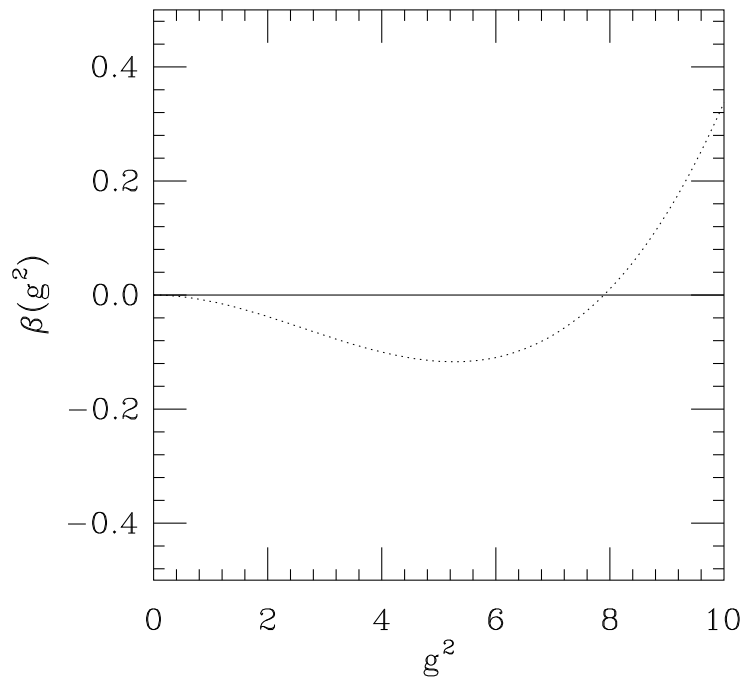
## Schrödinger functional

- Induce background field thru boundary conditions in  $L^4$  box
- Measure  $1/g^2(L)$  from response to small change in boundary conditions
- Extract mass anomalous dimension  $\gamma$  from scaling of pseudoscalar renormalization constant  $Z_P$ , on the same lattices

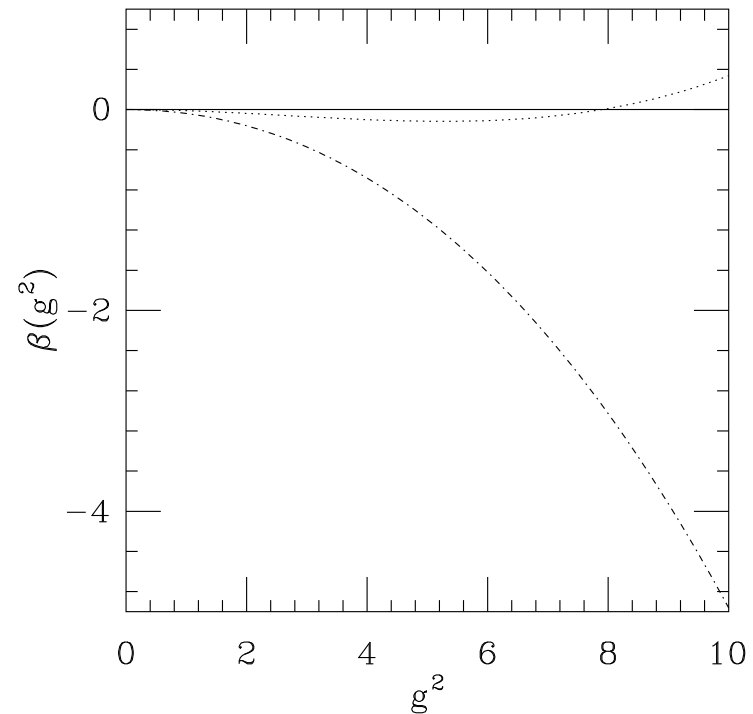
# Life inside (or near) the conformal window

Two-loop beta function,  $SU(2)$  with  $N_f = 2$  of:

adjoint



fundamental

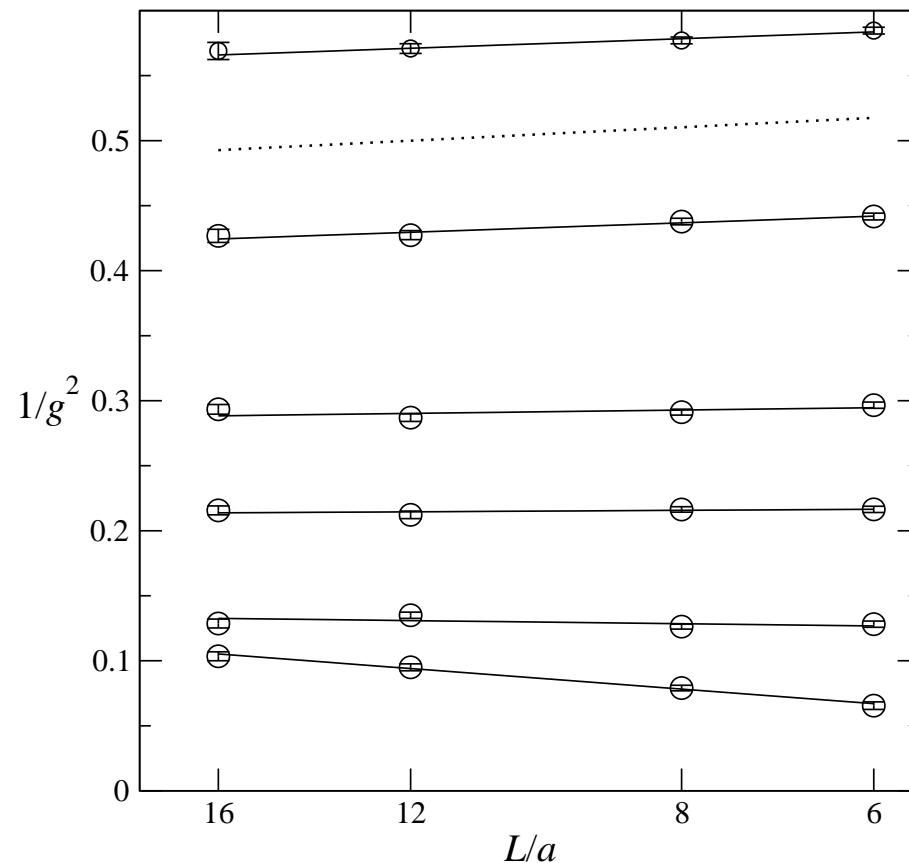


⇒ Strategy of ALPHA developed for QCD; performs poorly here

# Nearly constant beta function

define:  $\tilde{\beta}(1/g^2) \equiv \frac{d(1/g^2)}{d \log(L)} = 2\beta(g^2)/g^4$

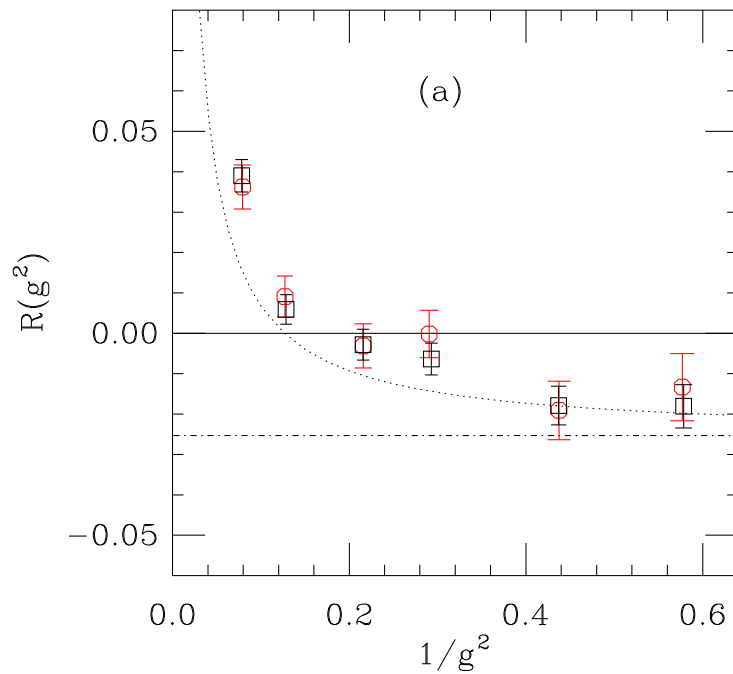
- Solution:  $1/g^2 \simeq a + b \log(L)$
- Exact at one loop
- Treat dataset at each bare  $\beta$  as separate fitting problem
- slope = beta fn.
- Systematics: add  $\log^2$  term, or remove smallest volume
- Slope changes sign  $\Rightarrow$  IRFP!



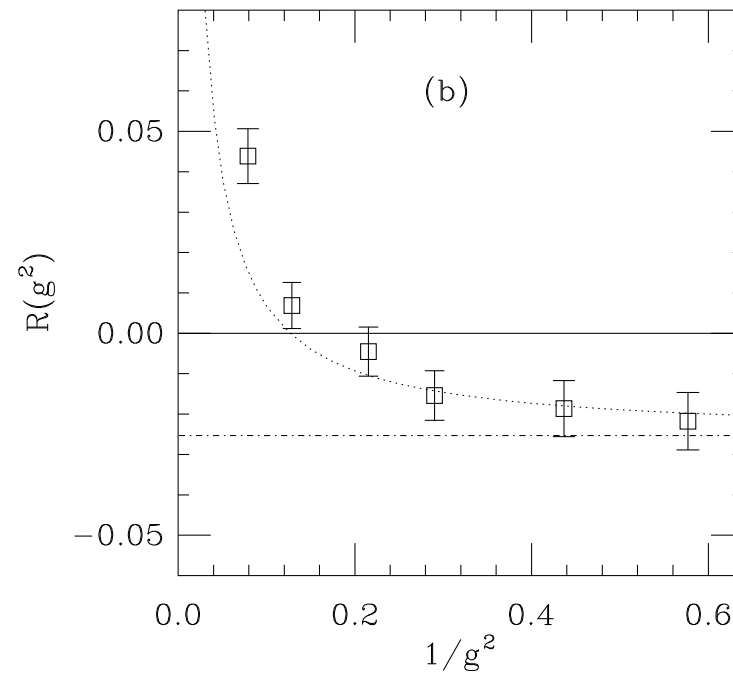
# SU(2)/adjoint

Fitting  $1/g^2$ :  $x = L/8$ , Black:  $L = 6, 8, 12, 16$ , Red:  $L = 8, 12, 16$

$$a + b \log(x)$$



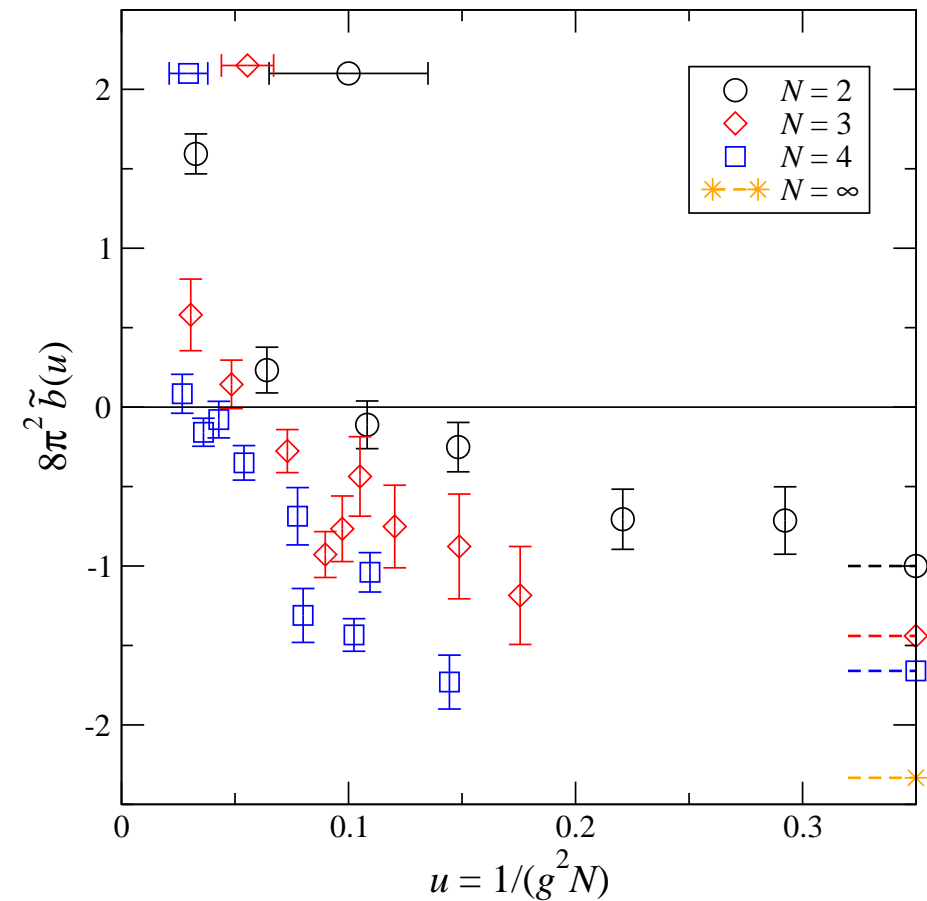
$$a + b \log(x) + c \log^2(x)$$



$$1/g_*^2 = 0.20(4)(3)$$

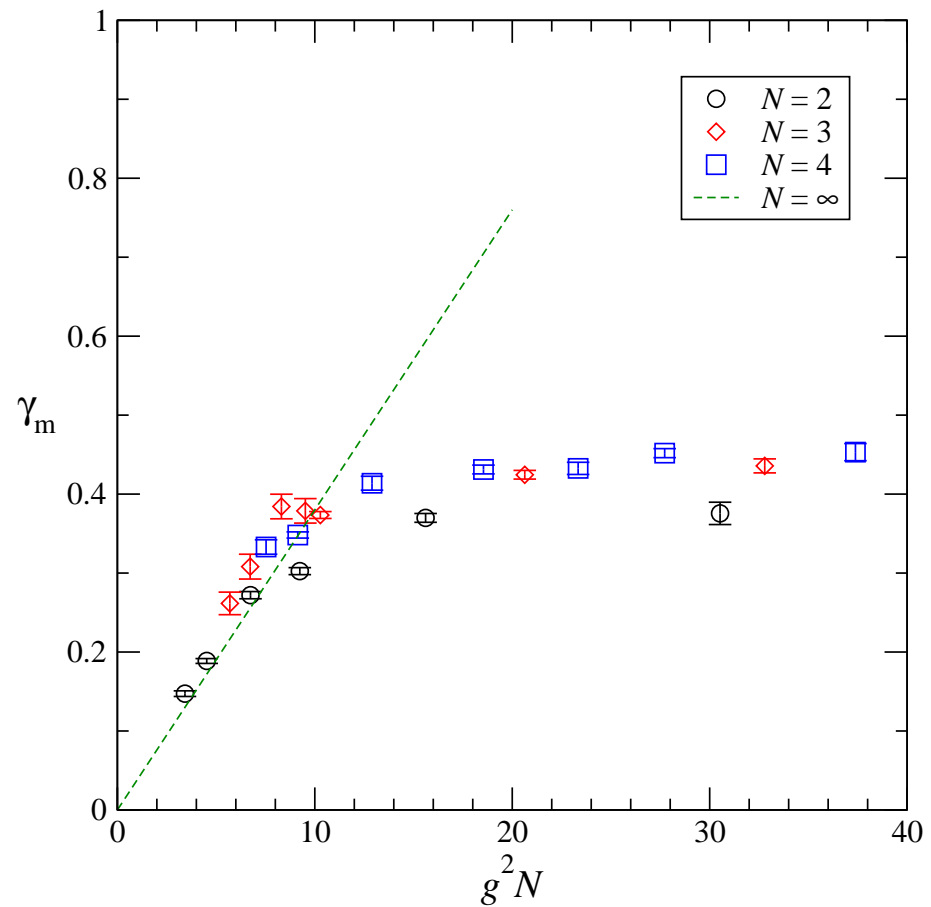
# SU(2)/adjoint, SU(3)/sextet, SU(4)/decuplet

- Qualitatively agrees with  $\beta$ [2-loop]
- Actual  $\beta > \beta$ [2-loop]
- With increasing  $N_c$ , IRFP moves to stronger coupling, or disappears
- Consistent with no IRFP of  $\beta$ [2-loop] for  $N_c \rightarrow \infty$  limit
- Can one use SU(4)/decuplet for walking technicolor?



# Mass anomalous dimension

- Dashed line:  
one-loop for  $N_c \rightarrow \infty$
- Remarkable universality
- Saturation:  $\gamma \lesssim 0.45$
- Missed by analytic calculations!
- No good for walking technicolor



## Part II. What nHYP links do for us

- Results from SU(3)/sextet

	$\beta$	$g_{SF}^2(L=6)$	$\kappa_c$	$m_c$	$M$
thin	6.0	$\sim 2$	0.1610	-0.89	1.9
fat	5.8	1.9	0.1283	-0.10	1.1

Critical hopping parameter  $\kappa_c = 1/(8 + 2m_c)$

Optimal Domain-Wall height  $M = 1 + |m_c|$  (to be explained later on)

- Similar improvement for optimal clover coefficient  
(big enough that we've decided to set  $c_{SW} = 1$ )
- Fat links allow us to probe much larger  $g^2(L)$

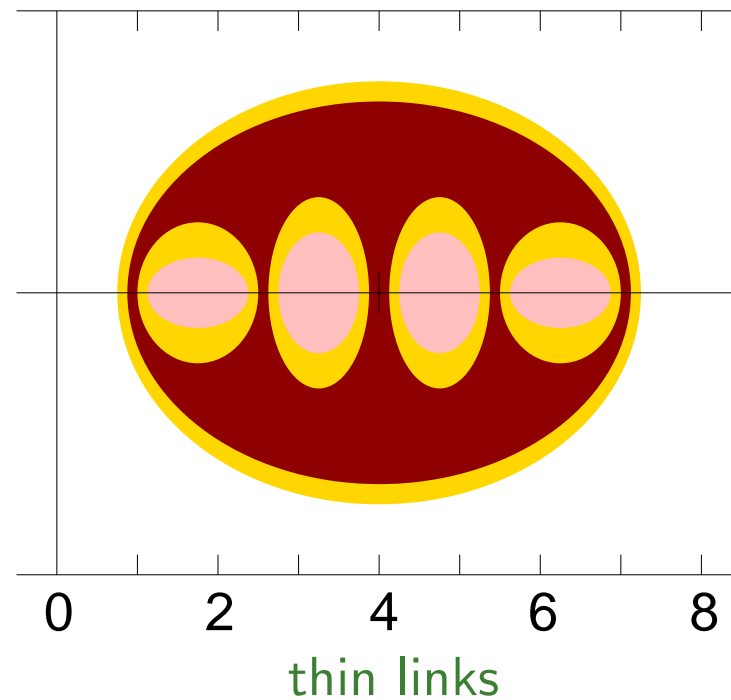
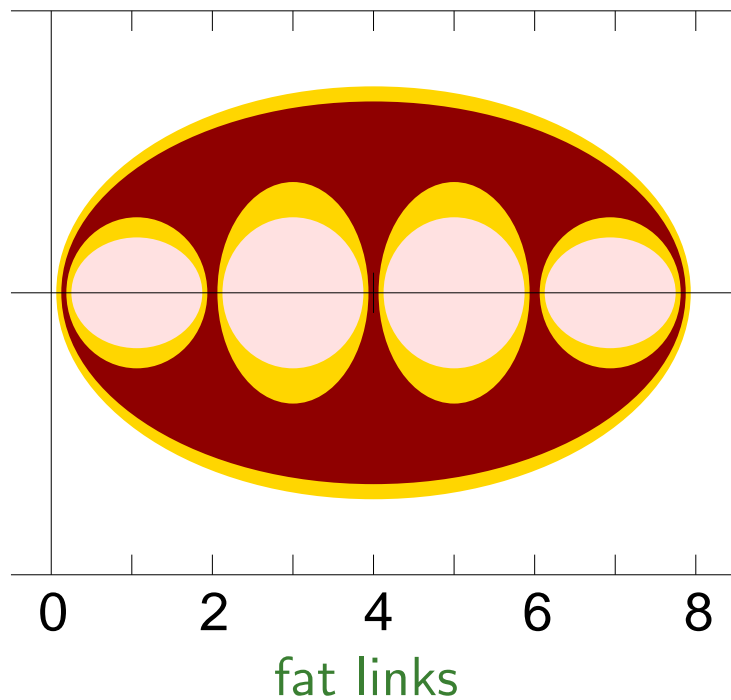


# DWF primer

- 5-d Wilson fermions, supercritical (negative bare mass), with chiral 4-d fields located near the boundaries  $s = 0$  and  $s = L$
- Effective 4-dim operator:  $D_{\text{eff}}(L) = 1 + \gamma_5 \tanh((L/2)a_5 H)$
- For  $L \rightarrow \infty$  it becomes a GW operator:  $D_{\text{eff}}(\infty) = 1 + \gamma_5 H / \sqrt{H^2}$
- DWF transfer matrix:  $T = \exp(-a_5 H) = 1 - a_5 H_W + \dots$   
where  $H_W = \gamma_5 D_W$  is the 4-d hermitian Wilson operator
- Penetration into the 5-d bulk from near-unity eigenvalues of  $T$ , hence from near-zero eigenvalues of  $H_W$
- PCAC:  $\partial_\mu A_\mu = 2m_q J_5 + 2m_{\text{res}}(J_5 + [\text{more}] \text{ lattice artifacts})$
- Need  $m_{\text{res}} \sim 10^{-3}$  (or smaller). Want  $L$  to be small.

## spectrum of $D_W$ in complex plane

- Optimal DW height = center of leftmost “void”
- Fat links tame the Wilson spectrum:  $m_c(\text{fat}) \ll m_c(\text{thin})$
- Expect many more ev’s inside the thin-links “void”:  $m_{\text{res}}(\text{fat}) \ll m_{\text{res}}(\text{thin})$



# Fat links and DWF

- DWF on staggered MILC lattices: [LHP Collaboration, . . . ]
  - thin:  $m_{\text{res}}$  “too big to measure”
  - fat:  $m_{\text{res}} \sim 10^{-3}$
- Locality
  - It’s true that  $D_W$  is less local with fat links
  - nHYP links only mildly nonlocal: [Hasenfratz, Hoffmann & Schaefer]  
Smeared link depends only on thin links that share a hypercube with it
- GW operators are never ultra-local
- Relevant notion: exponential localization of  $D_{\text{eff}}(L)$ .
  - Depending on suppression of near-zero eigenmodes, in principle  $D_{\text{eff}}(L)$  might be more local with fat links!

# Summary

- $SU(N)$ ,  $N = 2, 3, 4$ , with  $N_f = 2$  two-index sym rep Dirac fermions are **no good for walking technicolor** because mass anomalous dimension saturates at  $\gamma \lesssim 0.45$
- While several nice tricks are already being used to bring down  $m_{\text{res}}$  [improved gauge action, Dislocations Suppressing Determinant Ratio], and to speed up the inversion [Möbius], it's a good idea to **try DWF with fat links**

# Walking (extended) Technicolor

For acceptable flavor physics, candidate theories should have:

- Chiral symmetry breaking
- Small  $S$ -parameter
- Avoid massless GBs (except those eaten by  $W^\pm$  and  $Z$ )
- Condensate enhancement
  - $\Rightarrow$  need mass anomalous dimension  $\gamma \approx 1$
  - $\Rightarrow$  this is expected for nearly conformal (“walking”) theories
- $\Rightarrow$  Natural candidates:  $N_f = 2$  higher-irrep fermions

[Sannino *et al.*]

# Schrödinger functional

- Prescribe gauge field  $A_k(L, \eta)$  on time boundaries:  $t = 0$  and  $t = L$   
 $\Rightarrow$  induce background color-electric field in the bulk

$$\begin{aligned}\Gamma(L, \eta) \equiv -\log(Z) &= \text{tree-level} + \text{one-loop} + \dots \\ &= \left( \frac{1}{g_0^2(a)} + \frac{b_1}{32\pi^2} \log(L/a) + \dots \right) S_{cl}(\eta) \\ &= \frac{1}{g^2(L)} S_{cl}(\eta)\end{aligned}$$

- Obtain  $1/g^2(L)$  from variation w.r.t.  $\eta$ , which is an observable
- Extract mass anomalous dimension  $\gamma$  from scaling of pseudoscalar renormalization constant  $Z_P$ , on the same lattices

# Soft gauge action

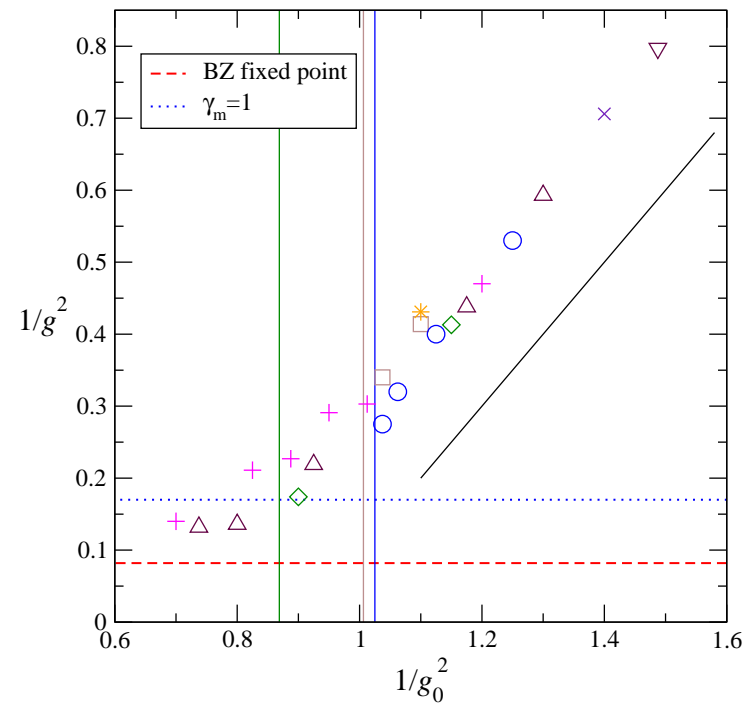
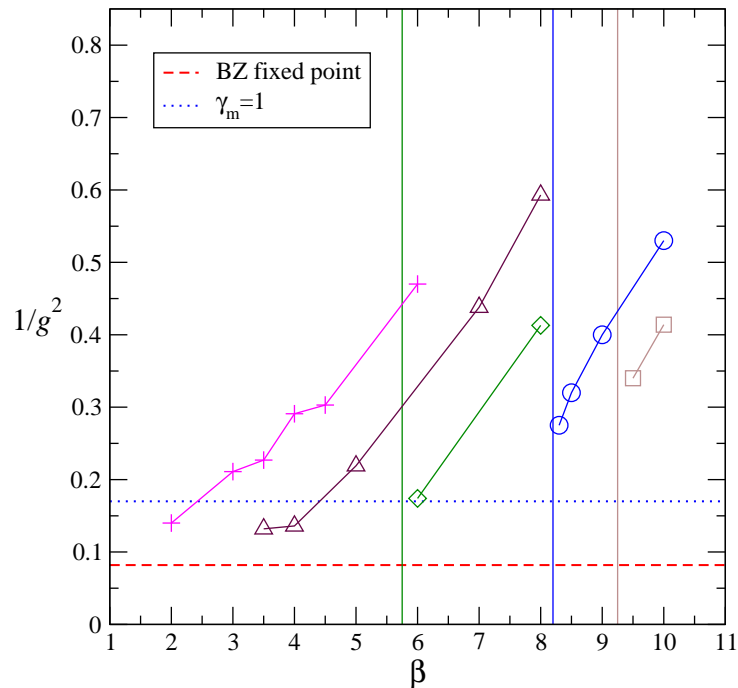
- Soft action:  $S_{\text{gauge}} = \beta \sum U_{\text{plaq}}(\text{thin links}) + \tilde{\beta} \sum U_{\text{plaq}}(\text{fat links})$
- Pushes 1st-order bulk transition back into stronger bare coupling
- SU(3)/sextet:
  - thin links: transition at  $g_{SF}^2 \approx 2.5$ , lattice artifacts for  $g_{SF}^2 \gtrsim 2.0$
  - fat links,  $\tilde{\beta} = 0$ : transition at  $g_{SF}^2 \approx 5.0$ , lattice artifacts for  $g_{SF}^2 \gtrsim 3.5$
  - fat links,  $\tilde{\beta} = 0.5$ : up to  $g_{SF}^2 \approx 11$  before running out of steam
- Stabilize nHYP reunitarization step:

$$V_\mu = \Omega_\mu (\Omega_\mu^\dagger \Omega_\mu)^{-1/2} \quad \Rightarrow \quad V_\mu = \Omega_\mu (\Omega_\mu^\dagger \Omega_\mu + \epsilon)^{-1/2}$$

We use  $\epsilon = 10^{-6}$

# Soft gauge action: weak-coupling universality

bare coupling:  $\frac{1}{g_0^2} = \frac{\beta}{N}T(f) + \frac{\tilde{\beta}}{d_R}T(R),$

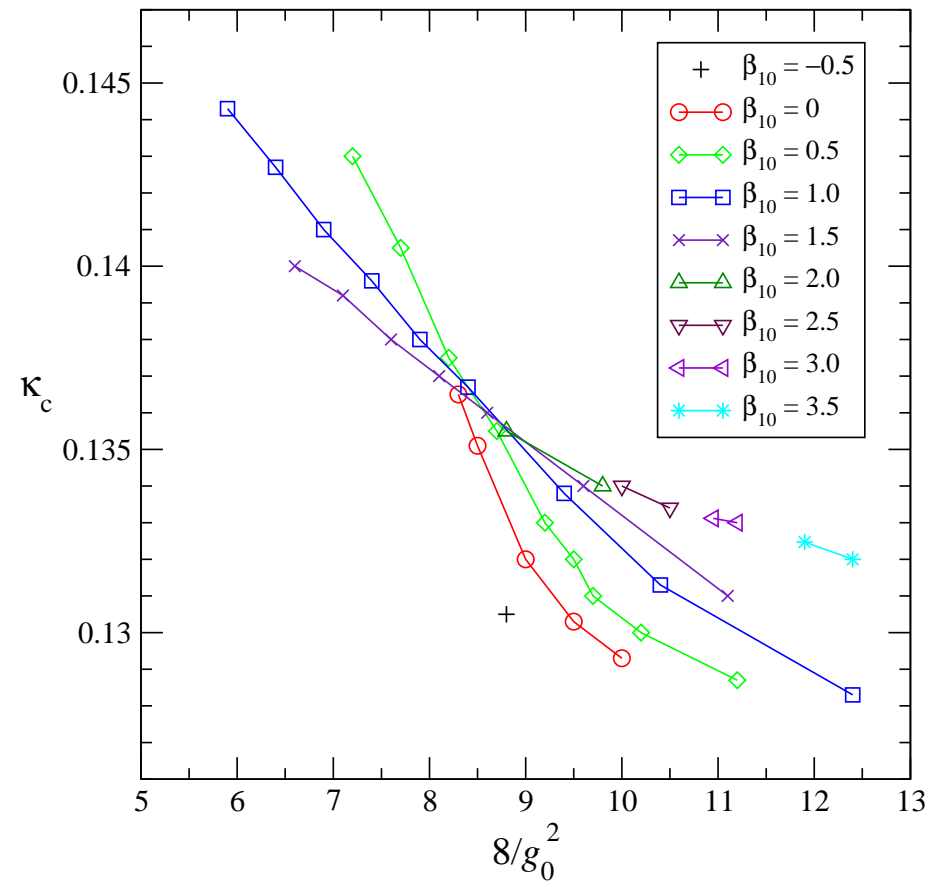


SU(4), Right to left:  $\tilde{\beta} = -0.5, 0.0, 0.5, 1.0, 1.5$

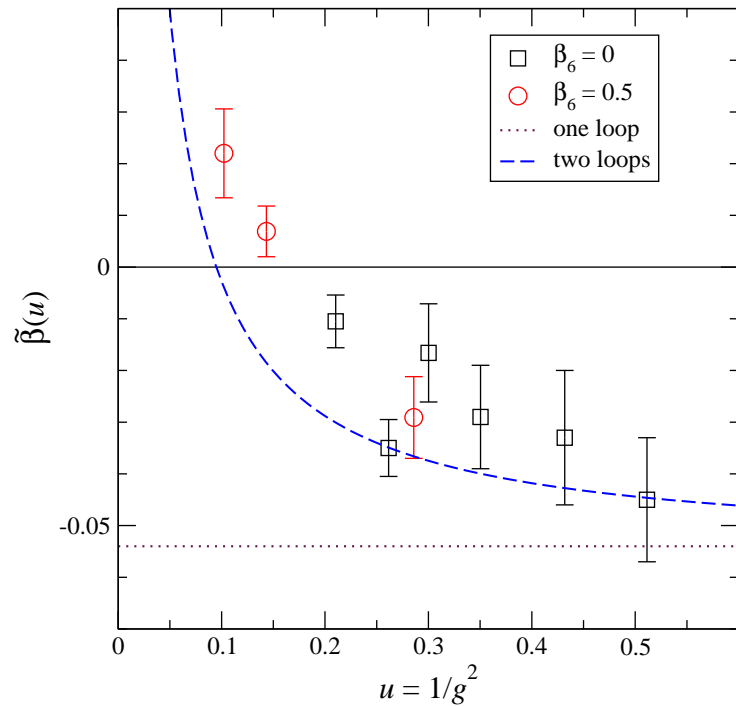


## Soft gauge action: $\kappa_c$

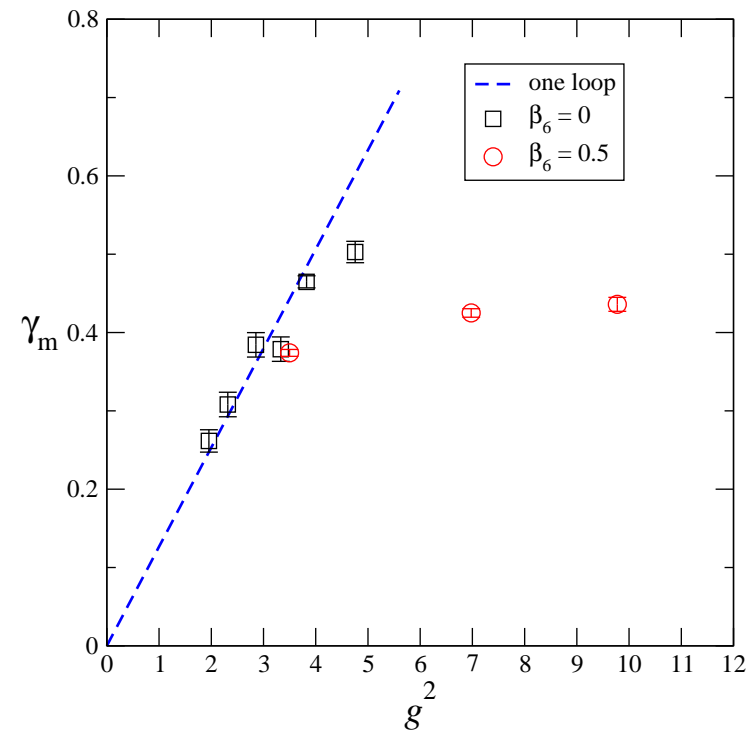
- Perturbation theory:  
 $\kappa_c$  increases with  $\tilde{\beta}$  at fixed  $1/g_0^2$
- Beyond  $8/g_0^2 \approx 9$  trend reverses



# Effects of the bulk transition, SU(3)



SF coupling: agreement



Mass anomalous dimension:  
discrepancy – due to proximity  
of bulk transition for  $\tilde{\beta} = 0$